Temperature-Dependent Orbital Degree of Freedom of a Bilayer Manganite by Magnetic Compton Scattering

Yinwan Li,1,2 P.A. Montano,1,3 J.F. Mitchell,2 B. Barbiellini,4 P.E. Mijnarends,4,5 S. Kaprzyk,4,6 and A. Bansil4

1Department of Physics, University of Illinois, Chicago, Illinois 60680, USA
2Materials Science Division, Argonne National Laboratory, Argonne, Illinois 60439, USA
4Physics Department, Northeastern University, Boston, Massachusetts 02115, USA
5Interfaculty Reactor Institute, Delft University of Technology, 2629 JB Delft, The Netherlands
6Academy of Mining and Metallurgy AGH, 30059 Kraków, Poland

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We have measured temperature-dependent magnetic Compton profiles (MCPs) from a single crystal of La12.8Sr0.2Mn2O7. The MCPs, which involved the scattering of circularly polarized x rays, are in general related to the momentum density of all the unpaired spins in the system. Nevertheless, we show that when the x-ray scattering vector lies along the [110] direction, the number of magnetic electrons of a specific symmetry, i.e., d electrons of \( \chi^2 - \gamma^2 \) symmetry, yield a distinct signature in the MCP, allowing us to monitor substantial changes in the occupancy of the \( d_{z^2-r^2} \) states over the investigated temperature range of 5–200 K. This study indicates that magnetic Compton scattering can provide a powerful window on the properties of specific magnetic electrons in complex materials.

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Layered manganites \([La_{1-x}Sr_xMnO_3]_2 SrO\) have been the subject of great current interest not only because they exhibit the colossal magnetoresistance effect [1,2], but also because these systems display a wide variety of magnetic properties [3] and undergo phase transitions associated with changes in the orbital degrees of freedom. Here we consider the application of the x-ray scattering spectroscopy in the deeply inelastic or so-called Compton scattering regime (i.e., involving a large transfer of energy and momentum in the scattering process) in order to gain a handle on the properties of magnetic electrons in the manganites.

The scattering cross section from magnetic electrons is typically several orders of magnitude smaller than in the more conventional charge scattering channel, and, therefore, magnetic Compton scattering experiments have become feasible on complex materials only in the last few years via the use of circularly polarized light at the high energy, high intensity synchrotron sources [4]. Previous magnetic Compton scattering work on the manganites, which has been limited to the double layer La manganite, has focused on the case where the x-ray scattering vector lies along the [100] or the [001] crystal direction [5–8]. The shape of the associated [100] and [001] magnetic Compton profiles (MCPs), however, is dominated by the \( d \) electrons of \( t_{2g} \) (i.e., \( xy, yz, \) and \( zx \)) symmetry. However, \( t_{2g} \) levels lie \( \sim 1 \) eV below the Fermi energy and are thus relatively “inert.” More central to the understanding of the manganites is the behavior of the magnetic states of \( e_g \) (i.e., \( x^2 - y^2 \) or \( 3z^2 - r^2 \)) symmetry.

This Letter focuses on the case where the x-ray scattering vector is chosen to lie along the [110] crystal direction. We show that the resulting [110] MCP, when Fourier transformed to real space, contains a remarkable signature of the number of \( d_{z^2-r^2} \) electrons through the presence of a well defined dip around 1 Å. Along the [110] direction, the shapes of various magnetic orbitals conspire in such a way as to render the depth of the aforementioned dip quite insensitive to the presence of other orbitals (\( t_{2g} \) and \( d_{x^2-y^2} \)). The [110] MCPs in \( La_{1.2}Sr_{1.8}Mn_2O_7 \) over the temperature range of 5–200 K are presented. In order to help interpret these experimental results, we also present the first all-electron computations of the MCP in the closely related compound \( LaSr_2Mn_2O_7 \). From our analysis of the temperature-dependent MCPs, we adduce that the number of \( d_{z^2-r^2} \) states increases from about 0.4 for each Mn atom at 5 K to around 0.5 at 160 K, but then decreases rapidly to the low-T value of \( \sim 0.4 \) by 200 K, suggesting that the system has undergone significant changes in the electronic structure. We emphasize that the 7 T magnetic field present during the experiments makes the samples electronically homogeneous and that this field is sufficiently strong to collapse polarons [9].

The magnetic Compton profile, \( J_{\text{mag}}(p_z) \), for momentum transfer along the scattering vector \( p_z \) is defined by

\[
J_{\text{mag}}(p_z) = J_1(p_z) - J_1(0),
\]

where \( J_1 (J_0) \) is the majority (minority) spin Compton profile. \( J_{\text{mag}} \) can be expressed in terms of a double integral of the spin density, \( \Delta \rho(p) \):

\[
J_{\text{mag}}(p_z) = \int \int \Delta \rho(p) dp_x dp_y,
\]

where \( \Delta \rho(p) \equiv \rho_1(p) - \rho_0(p) \).

MCP measurements were carried out using the elliptical multipole wiggler at beam line 11-B at the Advanced...
We have analyzed the origin of this peak extensively by considering a series of cuts through various high symmetry directions to establish that this feature arises from wave functions of $t_{2g}$ symmetry.

The interpretation of the momentum density and the MCPs is aided by Fig. 1, which shows two typical 2D projections (upper frames) of the magnetic momentum density $\Delta \rho(\mathbf{p})$, as well as two orthogonal cuts through the origin of $\Delta \rho(\mathbf{p})$ (lower frames). The (001) projection in Fig. 1(a), for instance, shows that the FS, seen by the presence of large and smaller squarish features, is essentially 2D in character. Underlying these FS features, the uncompensated magnetic contribution yields a strong peak at $(p_x, p_y) \sim (1.1, 1.1)$ a.u. in the projection of 1(a).

Indeed, the (001) and the (100) cuts in 1(c) and 1(d), respectively, both show a peak lying along the diagonal line as expected for an $xy$ or a $yz$ type orbital.

Recalling from Eq. (2) that the MCP involves a double integral or equivalently a 1D projection of $\Delta \rho(\mathbf{p})$, the form of the MCPs for momentum transfer along the principal symmetry directions can be understood straightforwardly from the 2D projections of Figs. 1(a) and 1(b). For example, the [100] MCP represents a further projection of the distribution of (a) on the horizontal [100] axis, so that it contains a $t_{2g}$ related peak around 1.1 a.u., which is also the case for the [001] MCP [as seen by projecting 1(b) on the vertical axis]. Such a peak is observed in our [100] and [001] MCPs and in those of Ref. [7]. These MCPs, however, are not shown here since our focus in this Letter is on the [110] MCP. We see from 1(b) that the region of high density is elongated along [110], so that the resulting [110] MCP (Fig. 2) is relatively flat unlike the [100] or the [001] MCPs, particularly after the experimental resolution is included. Figure 2 shows a remarkable level of agreement between theory and experiment, notwithstanding some discrepancies [17]. The structures in the unBroadened theory MCP (thin solid line) are due partly to the FS, but are washed out under the resolution of the present experiment (0.4 a.u.). Higher resolution measurements would be most interesting in gaining insight into FS signatures in the momentum density.

Figure 3 presents the temperature dependence of the [110] MCP. By invoking a molecular orbital picture [7], substantial changes observed in the shape of the [110] MCP with temperature can be understood. Although such an analysis is not rigorous, it nevertheless provides a handle on the important question of orbital occupancies. Assuming thus that the 3.6 Mn electrons essentially

Photon Source [10] on a high quality single crystal of La$_2$Sr$_{1.8}$Mn$_2$O$_7$ using circularly polarized photons. The sample is $10 \times 5 \times 2$ mm$^3$, where the shortest side is along the crystal $c$ axis. The sample was fixed to the holder with a specially made Al clip to avoid magnetic contamination. The MCPs were obtained using magnetic fields of 3 and 7 T along the three high symmetry directions [001], [100], and [110] for four different temperatures in each case: 5 K, which is well below the Curie temperature $T_c = 129$ K, and at 100, 160, and 200 K, the last temperature being well above $T_c$. The momentum resolution was $\sim 0.4$ a.u. $J_r(p_r)$ and $J_l(p_l)$ were measured by flipping the photon polarization. In this Letter, we focus on the [110] MCP and its temperature dependence under a 7 T field. We have also measured the MCP along [100] and [001] and observed no significant variation in the shape with temperature [5,6].

The electronic structure, magnetic momentum density, and the MCPs along principal symmetry directions were computed within an all-electron, fully charge and spin self-consistent semirelativistic Korringa-Kohn-Rostoker framework [11] for LaSr$_2$Mn$_2$O$_7$ in the $I4/mmm$ (No. 139) crystal structure [12]. All computations were carried out to a high degree of accuracy; e.g., the potentials were converged to better than $10^{-5}$ Ry. The momentum density was computed over a fine mesh of 129.6 $\times$ 10$^6$ points within a sphere of radius 14.9 a.u. in $\mathbf{p}$ space [13]. This data set was used to evaluate various spin-polarized projections and the MCPs along principal symmetry directions were obtained using magnetic fields of 3 and 7 T along the three high symmetry directions $\{001\}$, $\{100\}$ and $\{110\}$.

We focus on the $\{110\}$ MCP and its temperature dependence under a 7 T field. We have also measured the MCP along $\{100\}$ and $\{001\}$ and observed no significant variation in the shape with temperature [5,6].
The energy level diagram in Fig. 3 makes it clear that the occupancy of the $t_{2g}$ orbitals, which lie about 1 eV below $E_F$—well above $kT$—will change little with $T$, and that the variation with $T$ will occur mainly through the redistribution of 0.6 $e_g$ electrons between the $d_{x^2-y^2}$ and $d_{z^2-r^2}$ orbitals. For this purpose, we consider the reciprocal form factor, $B(r)$, which is defined as the Fourier transform of the spin momentum density 

$$B(r) = \int \Delta \rho(p) \exp(-ip \cdot r) dp.$$  

We now express $\Delta \rho(p)$ as a sum over the momentum density of the molecular orbitals $\psi_i^{MO}(p)$ of the magnetic electrons, weighted by their occupancies $n_i$:

$$\Delta \rho(p) = \sum_i n_i |\psi_i^{MO}(p)|^2.$$  

By transforming Eq. (4) into real space, it is easily shown that $B(r)$ gives the autocorrelation of the magnetic orbitals, $\psi_i^{MO}(r)$,

$$B(r) = \sum_i n_i \int \psi_i^{MO}(s)\psi_i^{MO}(s + r) ds,$$

where the integrals are the standard two-center overlap functions. Alternately, $B(r)$ along a given direction can be obtained directly by taking the 1D Fourier transform of the MCP along that direction by comparing Eqs. (2) and (3).

Figure 4 presents the $B(r)$ results along [110] obtained from the MCPs of Figs. 2 and 3. Notable features are the presence of a dip at $r = 1 \AA$ and oscillations at higher distances. As expected, the shape of the theory curve is in very good accord with the 5 K experimental data.
Systematic changes in shape with temperature can be interpreted by recalling that $B(r)$ gives the autocorrelation of the magnetic orbitals. We consider the dip in $B(r)$ first. We have computed the overlap integral in Eq. (5) for the $t_{2g}$ and $e_g$ orbitals using Slater-type orbitals (STO).

The key observation that emerges is that the large dip in $B(r)$ along [110] is produced essentially by the autocorrelation of the $d_{x^2-y^2}$ orbital. This can be understood qualitatively from the schematic picture of a $d_{x^2-y^2}$ orbital in Fig. 4. It is obvious that, when this orbital is translated along [110], the positive and negative lobes will overlap and yield a negative dip at a distance of the order of orbital dimensions. In contrast, the $d_{xy}$ and $d_{3z^2-r^2}$ orbitals give positive contributions at $r = 0.0025$; there are also small negative contributions from $d_{z^2}$ and $d_{3x^2-y^2}$ orbitals. These results indicate that the size of the negative dip in $B(r)$ is correlated with the $d_{x^2-y^2}$ occupancy. Computations in which all magnetic orbitals are included show that one can fit the minimum value, $B_{\text{min}}$, of $B(r)$ along [110] by a linear equation involving the occupancy $f$ of the $d_{x^2-y^2}$ orbital: $B_{\text{min}} = a f + b$, where $a = -0.130$ and $b = 0.004$ [20]. Using this fit, the $T$-dependent data of Fig. 4 give $f$ values of 0.38, 0.48, 0.53, and 0.38, corresponding to 5, 100, 160, and 200 K. Thus, at low $T$, the $d_{3z^2-r^2}$ occupancy is 0.22. By 160 K, this number decreases to 0.07, but then increases relatively rapidly to the low-$T$ value by 200 K. These variations correlate remarkably well with those in the apical oxygen distance, $D_{\text{apical}}$, in the MnO$_6$ octahedra. Reference [16] reports that $D_{\text{apical}}$ decreases up to $\approx 120$ K, but then begins to increase quite rapidly. A decrease in $D_{\text{apical}}$ is expected to raise the energy of the $d_{3z^2-r^2}$ levels and to reduce their occupancy or, equivalently, to increase the $d_{x^2-y^2}$ population, consistent with earlier observations.

Turning to the behavior of $B(r)$ beyond the dip, Fig. 4 shows that for $r \approx 1.5$ Å there is little overall change in shape up to 160 K. The 200 K curve, however, displays striking differences in that the broad feature around 2 Å shows that for $r > 120$ K, there is little overall change in the system lies near a multicritical point in the phase diagram, so that pinpointing the detailed nature of the aforementioned orbital change will generally be difficult.

In conclusion, our study shows that magnetic Compton scattering spectra taken under a field of 7 T allow access to the properties of the electronically homogeneous phase of the manganite over a wide range of temperatures. This strong magnetic field also collapses polarons. The shape of the spectrum for momentum transfer along the [110] direction contains a remarkable signature of the occupancy of the $d_{x^2-y^2}$ electrons. These results indicate that the Compton scattering can provide a powerful new spectroscopic window for investigating orbital physics and magnetic electrons in complex materials.

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[20] The formula for $B_{\text{min}}$ was obtained by choosing STOs and an occupancy of 3 for the $t_{2g}$ states, and of $f$ for the $d_{x^2-y^2}$ and $(0.6 - f)$ for the $d_{3z^2-r^2}$ states. The STOs yield $B(r) = (1 + t + c_2 f^2 + \cdots + c_6 f^6) \exp(-t)$, where $t = Z_{\text{eff}} r/3$, $Z_{\text{eff}}$ is the effective charge of the atomic potential, and the coefficients $c_i$ depend on $f$ [see, e.g., calculations of overlap integrals by M. P. Barnett, Int. J. Quantum Chem. 95, 791 (2003)]. In order to analyze changes in shape, all curves in Fig. 4 are normalized such that $B(0) = 1$. 